

where ϕ is another function.

To solve it we have .

$C = \text{constant}$ at const. temp

$\therefore d(c^2) = 0$ from eqn (A)

or, $d(u^2 + v^2 + w^2) = 0$

or, $u du + v dv + w dw = 0$ — (2)

Now taking differential of eqn (1) both sides, we have

$$f'(u) du f(v) f(w) + f'(v) dv f(u) f(w) + f'(w) dw f(u) f(v) = 0$$

where dashes represents 1st derivative

Dividing above eqn by $f(u) f(v) f(w)$

We have

$$\frac{f'(u)}{f(u)} du + \frac{f'(v)}{f(v)} dv + \frac{f'(w)}{f(w)} dw = 0$$
 — (3)

Multiplying eqn (2) by constant β both sides we get —

$$\beta u du + \beta v dv + \beta w dw = 0$$
 (3i)

Now combining (3) & (3i) by Laplace's method of undetermined multipliers

We have

$$\left(\frac{f'(u)}{f(u)} + \beta u \right) du + \left(\frac{f'(v)}{f(v)} + \beta v \right) dv + \left(\frac{f'(w)}{f(w)} + \beta w \right) dw = 0$$
 — (4)

identity
 $\frac{d}{dx}(f(x)g(x)) = f'(x)g(x) + f(x)g'(x)$

Now this expression is an identity.
 Hence separate terms must vanish.

Taking 1st term we get

$$\left(\frac{f'(u)}{f(u)} + \beta u \right) du = 0$$

$$\text{or, } \frac{f'(u)}{f(u)} du = -\beta u du$$

Integrating we get

$$\log f(u) = -\beta \frac{u^2}{2} + \text{const.}$$

$$= -\beta \frac{u^2}{2} \log_e e + \log a$$

$$= \log_e e^{-\beta \frac{u^2}{2}} + \log a$$

$$= \log_e e^{-bu^2} + \log a$$

$$\text{where } b = \beta/2$$

$$\text{or } \log f(u) = \log a e^{-bu^2}$$

$$\text{Similarly } \text{or } f(u) = a e^{-bu^2}$$

$$\text{Similarly } f(v) = a e^{-bv^2}$$

$$\text{and } f(w) = a e^{-bw^2}$$

These function are determined and we have

$$f(u)f(v)f(w) = a^3 e^{-b(u^2+v^2+w^2)}$$

And the number of molecules lying between velocity range u and $u+du$, v and $v+dv$, w & $w+dw$ in elementary volume is given by

where
 $\log_e e = 1$
 $a = \text{const.}$
 $\log_e 1 = 0$
 $\log_a b = \frac{1}{a} \log b$

$$dn = n a^3 e^{-b(u^2+v^2+w^2)} du dv dw$$

The constant $a = \sqrt{\frac{b}{\pi}}$ and $b = \frac{m}{2kT}$

where $k =$ Boltzmann constant.

$T =$ absolute temp.

$m =$ mass of a molecule

Hence

$$dn = n \left(\frac{m}{2\pi kT} \right)^{3/2} e^{-\frac{m}{2} \left(\frac{u^2+v^2+w^2}{kT} \right)} du dv dw$$

(5)

We can now obtain the number of molecules lying between c & $c+dc$.

Where we have the volume $du dv dw$ is velocity space and in place of this elementary volume. The volume element in spherical polar co-ordinates (r, θ, ϕ)

$$= dc \cdot c d\theta \cdot c \sin\theta d\phi$$

and may be substituting for $du dv dw$ in eqn (5) and velocities are now between c & $c+dc$, θ & $\theta+d\theta$, ϕ & $\phi+d\phi$.

To find the total no. of molecules between c & $c+dc$, we have to integrate the expression for all possible values of ϕ & θ .

Hence

$$dn_c = n a^3 e^{-bc^2} c^2 dc \int_0^{2\pi} d\phi \int_0^{\pi} \sin\theta d\theta$$

$$dn_c = 4\pi n a^3 e^{-bc^2} c^2 dc$$

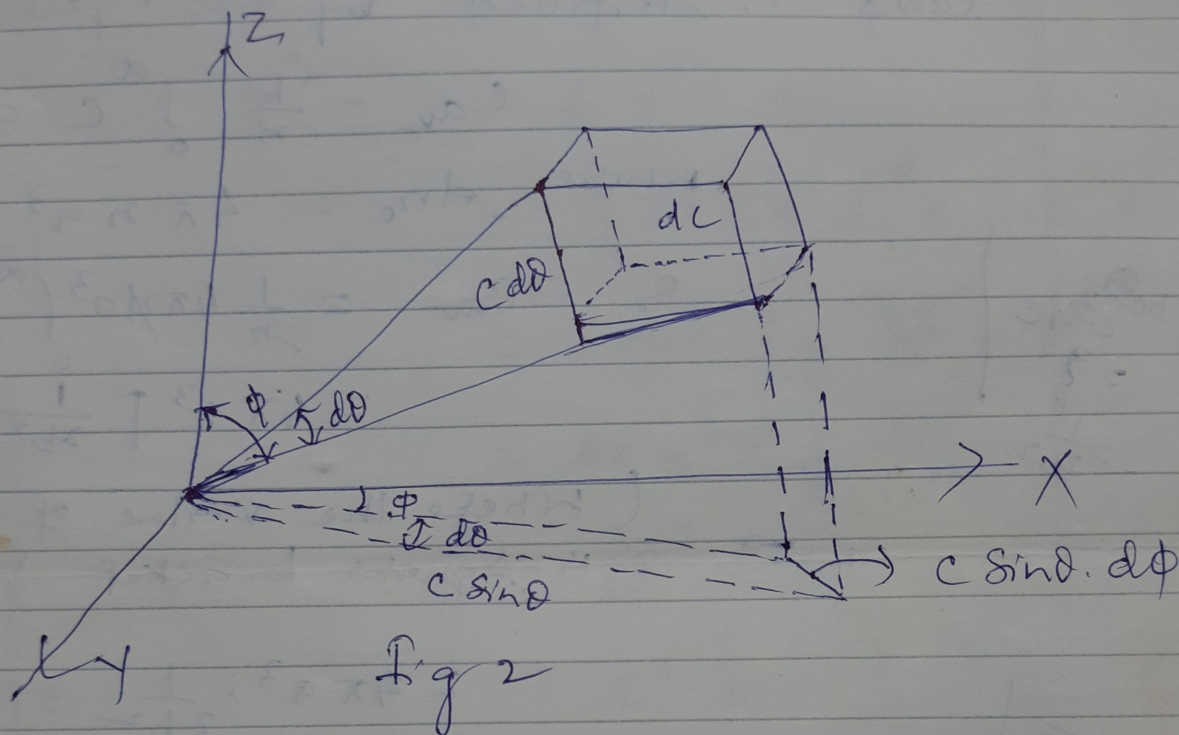
(6)

which is the distribution law of Maxwell.

$[\phi]_0^{2\pi}$
 $[-0 = 2\pi]$
 $d\theta = [-\cos\theta]_0^{\pi}$
 $[\pi - \cos\theta]$
 $[-1 - 1]$
 $[-2]$
 $\times 2$

Which is the distribution law of Maxwell
 slow putting $bc^2 = x^2$
 we have

$$dn_c = 4\pi \pi^{-1/2} x^2 e^{-x^2} dx$$



If we plot the function $y = 4\pi^{-1/2} e^{-x^2}$ against x , we get the Maxwellian curves as shown in fig below.

And the number of molecules dn_c whose speed lies between u & $u+du$ is proportional to shadowed area -

